

ON ASYMPTOTIC BEHAVIOR OF SOLUTION OF ONE NONLINEAR
ONE-DIMENSIONAL INTEGRO-DIFFERENTIAL ANALOGUE OF MAXWELL'S
SYSTEM

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Abstract. Large time behavior of solutions of a system of nonlinear integro-differential equations associated with the penetration of a magnetic field into a substance is studied. Initial-boundary value problem with Dirichlet boundary conditions is considered. Exponential stabilization of solution is established. Corresponding finite difference scheme is considered as well.

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A great variety of applied problems are modeled by the systems of nonlinear integro-differential equations. Such systems arise for instance for mathematical modeling of the process of penetrating of magnetic field in the substance. Corresponding Maxwell's system [1] can be rewritten in the following form [2]:

$$\frac{\partial H}{\partial t} = -rot \left[a \left(\int_0^t |rot H|^2 d\tau \right) rot H \right], \quad (1)$$

where $H = (H_1, H_2, H_3)$ is a vector of the magnetic field and function $a = a(S)$ is defined for $S \in [0, \infty)$.

If the magnetic field has the form $H = (0, U, V)$ and $U = U(x, t)$, $V = V(x, t)$, then we obtain the following system of nonlinear integro-differential equations:

$$\begin{aligned} \frac{\partial U}{\partial t} &= \frac{\partial}{\partial x} \left[a \left(\int_0^t \left[\left(\frac{\partial U}{\partial x} \right)^2 + \left(\frac{\partial V}{\partial x} \right)^2 \right] d\tau \right) \frac{\partial U}{\partial x} \right], \\ \frac{\partial V}{\partial t} &= \frac{\partial}{\partial x} \left[a \left(\int_0^t \left[\left(\frac{\partial U}{\partial x} \right)^2 + \left(\frac{\partial V}{\partial x} \right)^2 \right] d\tau \right) \frac{\partial V}{\partial x} \right]. \end{aligned} \quad (2)$$

Many scientific works are devoted to uniqueness and solvability of various problems for (2) type models (see, for example, [2]-[8] and references therein). The existence of the global solutions for initial-boundary value problems of such models have been proved in [2]-[5],[8] by using the modified Galerkin and compactness methods [9],[10]. For solvability and uniqueness properties of such type models see also [6],[7] and number

of other works as well. In [6] some generalization of equations of type (1) is proposed. In this case analogous of the model type (2) has the following form:

$$\begin{aligned} \frac{\partial U}{\partial t} &= a \left(\int_0^t \int_0^1 \left[\left(\frac{\partial U}{\partial x} \right)^2 + \left(\frac{\partial V}{\partial x} \right)^2 \right] dx d\tau \right) \frac{\partial^2 U}{\partial x^2}, \\ \frac{\partial V}{\partial t} &= a \left(\int_0^t \int_0^1 \left[\left(\frac{\partial U}{\partial x} \right)^2 + \left(\frac{\partial V}{\partial x} \right)^2 \right] dx d\tau \right) \frac{\partial^2 V}{\partial x^2}. \end{aligned} \quad (3)$$

The asymptotic behavior of the solutions of (2) and (3) type models have been object of intensive research (see, for example, [8],[11]-[19]).

The purpose of this note is to study the asymptotic behavior of solutions of the initial-boundary value problem for the system (3) in the case $a(S) = 1 + S$ and confirmation of theoretical results by numerical experiments.

In the domain $[0, 1] \times [0, \infty)$ for the system (3) with $a(S) = 1 + S$ let us consider the following initial-boundary value problem:

$$\begin{aligned} \frac{\partial U}{\partial t} &= \left\{ 1 + \int_0^t \int_0^1 \left[\left(\frac{\partial U}{\partial x} \right)^2 + \left(\frac{\partial V}{\partial x} \right)^2 \right] dx d\tau \right\} \frac{\partial^2 U}{\partial x^2}, \\ \frac{\partial V}{\partial t} &= \left\{ 1 + \int_0^t \int_0^1 \left[\left(\frac{\partial U}{\partial x} \right)^2 + \left(\frac{\partial V}{\partial x} \right)^2 \right] dx d\tau \right\} \frac{\partial^2 V}{\partial x^2}, \end{aligned} \quad (4)$$

$$U(0, t) = V(0, t) = 0, \quad U(1, t) = \psi_1, \quad V(1, t) = \psi_2, \quad (5)$$

$$U(x, 0) = U_0(x), \quad V(x, 0) = V_0(x), \quad (6)$$

where $\psi_1 = Const \geq 0$, $\psi_2 = Const \geq 0$; $U_0 = U_0(x)$ and $V_0 = V_0(x)$ are given functions.

The existence and uniqueness of the solution of such type problems for one equation in suitable classes are proved in [8]. One must note that asymptotic behavior of solution for problem (4)-(6) with homogeneous as well as non-homogenous boundary conditions is studied in [17]. For problem with homogeneous boundary condition exponential stabilization while for problem with nonhomogeneous boundary condition on one side of lateral boundary power-like stabilization is obtained.

We use usual Sobolev spaces $H^k(0, 1)$ and $H_0^k(0, 1)$.

Theorem 1. *If $U_0, V_0 \in H_0^1(0, 1)$, then for the solution of problem (4)-(6) the following estimate is true*

$$\|U\| + \left\| \frac{\partial U}{\partial x} \right\| + \|V\| + \left\| \frac{\partial V}{\partial x} \right\| \leq C \exp\left(-\frac{t}{2}\right).$$

Note that here and below C denotes positive constants independent of t and $\|\cdot\|$ denotes norm in the space $L_2(0, 1)$.

Note that Theorem 1 gives exponential stabilization of the solution of the problem (4)-(6) in the norm of the space $H^1(0, 1)$. The stabilization is also achieved in the norm of the space $C^1(0, 1)$. In particular, the following statement takes place.

Theorem 2. *If $U_0, V_0 \in H^3(0, 1) \cap H_0^1(0, 1)$, then for the solution of problem (4)-(6) the following relations hold:*

$$\left| \frac{\partial U(x, t)}{\partial x} \right| \leq C \exp\left(-\frac{\alpha t}{2}\right), \quad \left| \frac{\partial V(x, t)}{\partial x} \right| \leq C \exp\left(-\frac{\alpha t}{2}\right),$$

$$\left| \frac{\partial U(x, t)}{\partial t} \right| \leq C \exp\left(-\frac{\beta t}{2}\right), \quad \left| \frac{\partial V(x, t)}{\partial t} \right| \leq C \exp\left(-\frac{\beta t}{2}\right),$$

where $0 < \alpha = \text{Const} < 1$, $0 < \beta = \text{Const} < \alpha$.

Following lemma is necessary to prove Theorem 2.

Lemma 1. *For the solution of problem (4)-(6) the following estimate is true*

$$\left\| \frac{\partial U(x, t)}{\partial t} \right\| + \left\| \frac{\partial V(x, t)}{\partial t} \right\| \leq C \exp\left(-\frac{t}{2}\right).$$

In the rectangle $[0, 1] \times [0, T]$, where T is a positive constant, we consider problem (4)-(6). We assume that $U_0 = U_0(x)$ and $V_0 = V_0(x)$ are sufficiently smooth given functions of their arguments.

Using usual notations [20] we correspond to the problem (4)-(6) the difference scheme:

$$\frac{u_i^{j+1} - u_i^j}{\tau} = \left\{ 1 + \tau h \sum_{l=1}^M \sum_{k=1}^{j+1} [(u_{\bar{x}, l}^k)^2 + (v_{\bar{x}, l}^k)^2] \right\} u_{\bar{x}, i}^{j+1},$$

$$\frac{v_i^{j+1} - v_i^j}{\tau} = \left\{ 1 + \tau h \sum_{l=1}^M \sum_{k=1}^{j+1} [(u_{\bar{x}, l}^k)^2 + (v_{\bar{x}, l}^k)^2] \right\} v_{\bar{x}, i}^{j+1},$$

$$i = 1, 2, \dots, M-1; \quad j = 0, 1, \dots, N-1, \quad (7)$$

$$u_0^j = v_0^j = 0, \quad u_M^j = \psi_1, \quad v_M^j = \psi_2, \quad j = 0, 1, \dots, N, \quad (8)$$

$$u_i^0 = U_{0,i}, \quad v_i^0 = V_{0,i} \quad i = 0, 1, \dots, M. \quad (9)$$

Many scientific works are devoted to the construction of discrete analogues for (2) and (3) type models (see, for example, [14],[17],[18],[21],[22] and references therein).

It is not difficult to obtain for (7)-(9) the following estimation:

$$\|u^n\|_h^2 + \sum_{j=1}^n \|u_{\bar{x}}^j\|_h^2 \tau \leq C, \quad \|v^n\|_h^2 + \sum_{j=1}^n \|v_{\bar{x}}^j\|_h^2 \tau \leq C, \quad n = 1, 2, \dots, N. \quad (10)$$

The a-priori estimate (10) guarantees the stability of the scheme (7)-(9).

Theorem 3. *If the problem (4)-(6) has a sufficiently smooth unique solution $U = U(x, t)$, $V = V(x, t)$, then exists the unique solution $u^j = (u_1^j, u_2^j, \dots, u_{M-1}^j)$, $v^j = (v_1^j, v_2^j, \dots, v_{M-1}^j)$, $j = 1, 2, \dots, N$ of the finite difference scheme (7)-(9) which tends to the $U^j = (U_1^j, U_2^j, \dots, U_{M-1}^j)$, $V^j = (V_1^j, V_2^j, \dots, V_{M-1}^j)$ for $j = 1, 2, \dots, N$ as $\tau \rightarrow 0$, $h \rightarrow 0$ and the following estimates are true*

$$\|u^j - U^j\|_h \leq C(\tau + h), \quad \|v^j - V^j\|_h \leq C(\tau + h), \quad j = 1, 2, \dots, N.$$

We now comment on the numerical implementation of the discrete problem (7)-(9). Note that (7) can be rewritten as:

$$\begin{aligned} \frac{u_i^{j+1} - u_i^j}{\tau} - A^{j+1} \frac{u_{i+1}^{j+1} - 2u_i^{j+1} + u_{i-1}^{j+1}}{h^2} &= 0, \\ \frac{v_i^{j+1} - v_i^j}{\tau} - A^{j+1} \frac{v_{i+1}^{j+1} - 2v_i^{j+1} + v_{i-1}^{j+1}}{h^2} &= 0, \\ i = 1, 2, \dots, M-1, \quad j = 0, 1, \dots, N-1, \end{aligned}$$

where

$$A^j = 1 + \tau h \sum_{\ell=1}^M \sum_{k=1}^j \left[\left(\frac{u_\ell^k - u_{\ell-1}^k}{h} \right)^2 + \left(\frac{v_\ell^k - v_{\ell-1}^k}{h} \right)^2 \right].$$

In order to rewrite this in matrix form, we define the vectors $\mathbf{u}^j = [u_1^j, u_2^j, \dots, u_{M-1}^j]^T$ and similarly \mathbf{v}^j . We also define the symmetric tridiagonal $(M-1) \times (M-1)$ matrix \mathbf{T} as follows

$$\mathbf{T}_{rs}^{j+1} = \begin{cases} -\frac{1}{h^2} A^{j+1}, & s = r-1, \\ \frac{2}{h^2} A^{j+1}, & s = r, \\ -\frac{1}{h^2} A^{j+1}, & s = r+1, \\ 0, & \text{otherwise.} \end{cases}$$

Thus the system (7) becomes

$$\frac{1}{\tau} \begin{bmatrix} \mathbf{u}^{j+1} \\ \mathbf{v}^{j+1} \end{bmatrix} - \frac{1}{\tau} \begin{bmatrix} \mathbf{u}^j \\ \mathbf{v}^j \end{bmatrix} + \begin{bmatrix} \mathbf{T}^{j+1} & 0 \\ 0 & \mathbf{T}^{j+1} \end{bmatrix} \begin{bmatrix} \mathbf{u}^{j+1} \\ \mathbf{v}^{j+1} \end{bmatrix} = 0. \quad (11)$$

We use Newton's method to solve the nonlinear system (11). Let $\mathbf{P}^j = [\mathbf{u}^j, \mathbf{v}^j]^T$ and define

$$\mathbf{H}(\mathbf{P}^{j+1}) = \frac{1}{\tau} \mathbf{P}^{j+1} - \frac{1}{\tau} \mathbf{P}^j + \hat{\mathbf{T}}^{j+1} \mathbf{P}^{j+1}, \quad (12)$$

where $\hat{\mathbf{T}}^{j+1}$ is the 2 by 2 block diagonal matrix with \mathbf{T}^{j+1} on diagonal. Newton's method for the system (12) is given by

$$\nabla \mathbf{H}(\mathbf{P}^{j+1}) \Big|^{(n)} \left(\mathbf{P}^{j+1} \Big|^{(n+1)} - \mathbf{P}^{j+1} \Big|^{(n)} \right) = -\mathbf{H}(\mathbf{P}^{j+1}) \Big|^{(n)}.$$

It is well known that if H_i are three times continuously differentiable in a region containing the solution and the Jacobian does not vanish in that region, then Newton's method converges at least quadratically (see, for example, [23]). In our case the Jacobian is the matrix ∇H in which the term $\frac{1}{\tau}$ on diagonal ensures that it doesn't vanish. The differentiability is guaranteed, since ∇H is quadratic. So, we obtain convergence of the considered iterative method.

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