

ON SOME NEW REPRESENTATIONS OF HOLOMORPHIC
FUNCTIONS IN LATTICED DOMAINS

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Abstract

In this paper some new representations of holomorphic functions in latticed domains are obtained by means of conformal mapping singular integral equation method.

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Various problems of Mathematical Physics are reduced to the Dirichlet planar problem for the Laplace equation in a latticed domains. It is well known that this problem is closely connected with the boundary value problems for the holomorphic functions. For example the theory of hydro-turbines, the 3D motion of particles in a torus, the wave propagation in the sells (the size of sell is so small, that the whole domain we can consider as the infinite) [2,6,7,9].

At first, let us recall some notations from the theory of doubly-periodic functions [1, 3, 4, 6].

Consider a complex z -plane (\mathbb{C}), and two numbers $\omega_1, \omega_2 > 0$.

Definition 1 A set $D \subset \mathbb{C}$ is called doubly-periodic set if $z \in \mathbb{C}$ implies $z + 2m\omega_1 + 2ni\omega_2 \in D$, $m, n = 0, \pm 1, \dots$ points z and $z + 2m\omega_1 + 2ni\omega_2$ are called the congruent points.

Definition 2 The interior of the parallelogram with the vertexes $0, 2\omega_1, 2\omega_1 + 2i\omega_2, 2i\omega_2$ and with sides $[0\omega_1]$ and $[0, 2i\omega_2]$ is called the fundamental parallelogram [1,3,4, 6]. The interior domain of the fundamental parallelogram we denote by S_{00} .

On the figures 1 and 2 examples of doubly-periodic sets are given.

We also use definitions introduced in [5].

Definition 3 The function $\Phi(z)$ defined in the doubly-periodic domain is called exponentially doubly-quasi periodic if the following conditions are fulfilled

$$\Phi(z + 2\omega_1) = \Phi(z) \exp(P_{k_1}(z)), \quad \Phi(z + 2i\omega_2) = \Phi(z) \exp(Q_{k_2}(z)),$$

where P_{k_1} and Q_{k_2} are the definite polynomials of the k_1 and k_2 orders respectively.

This class of functions we denote by $P_e(k), k = \max(k_1, k_2)$ [5].

Example. The Weierstrass “sigma-function” $\sigma(z)$ is of the class $P_e(1)$. Indeed

$$\sigma(z+2\omega_1) = \sigma(z) \exp(\delta_1(z+\omega_1)+\pi i), \quad \sigma(z+2i\omega_2) = \sigma(z) \exp(\delta_2(z+i\omega_2)+\pi i),$$

where δ_1 and δ_2 are the definite constants [1,4,6].

Definition 4 A function $F(z)$ defined in the doubly-periodic domain D is called polynomially doubly quasi-periodic of k -order with the periods $2\omega_1$ and $2i\omega_2$ if the following conditions are fulfilled

$$F(z+2\omega_1) = F(z) + P_{k_1}(z), \quad F(z+2i\omega_2) = F(z) + Q_{k_2}(z), \quad z \in D,$$

where P_{k_1}, Q_{k_2} , are the definite polynomials of degree k_1 and k_2 respectively. This polynomials we call the proper polynomials of the function $F(z)$.

This class of functions we denote by $\mathcal{P}(k), k = \max(k_1, k_2)$ [5]. The class $\mathcal{P}(0)$ is the class of doubly quasi-periodic functions.

This class of functions we denote by $P_e(k), k = \max(k_1, k_2)$ [5].

Example. The Weierstrass “zeta-function” $\zeta(z)$ is doubly quasi-periodic function [1,4,6]

$$\zeta(z + 2\omega_1) = \zeta(z) + \delta_1, \quad \zeta(z + 2i\omega_2) = \zeta(z) + \delta_2.$$

In a complex z -plane let us consider the doubly-periodic line L which is a union of a countable number of smooth non-intersected contours $L_{mn}^j, j = 1, 2, \dots, k, m, n = 0, \pm 1, \dots$ doubly-periodically distributed with periods $2\omega_1$ and $2i\omega_2$ in the whole z -plane

$$L = \bigcup_{m,n=-\infty}^{\infty} L_{mn},$$

$$L_{mn} = \bigcup_{j=1}^k L_{mn}^j, \quad L_{mn}^{j_1} \cap L_{mn}^{j_2} = \emptyset, \quad j_1 \neq j_2.$$

In the sequel we will consider the case, when the contours L_{mn} are closed.

By S^+ we denote an infinite region bounded by the contour L . The positive direction on L will be taken such that S^+ remains on the left.

The union of the finite domains S_{mn}^- contained in every $L_{mn}, m, n = 0, \pm 1, \dots$ respectively will be denoted by S^- .

We consider the case when L_{00} is simple closed contour symmetric with respect to the axis $x = \omega_1$ and $y = i\omega_2$. Let us consider the following problem

Problem 1 To find the real function $u(x, y), u(x, y) \geq 0$, sub-harmonic in S^+ , continuous in $S^- + L$, having simple zeros at the finite number of points $b_i, i = 1, 2, \dots, l$ of the area S_{00} and satisfying the boundary condition

$$u = f(t), \quad t \in L, \quad (1)$$

where $f(t), f(t) > 0$ is real function given on L satisfies the H condition on L_{00} , of $P_e(1)$ class and symmetric with respect to the axis $y = i\omega_2$. $u(x, y)$ is also assumed to be symmetric and to belong to the class $P_e(1)$. We admit that proper polynomials of $f(t)$ are

$$\begin{aligned} P(z) &= \operatorname{Re} \alpha_1 z + \beta_1, \\ Q(z) &= \operatorname{Re} \alpha_2 z + \beta_2, \end{aligned}$$

$\alpha_1, \beta_1, \alpha_2, \beta_2$ are given constants.

With this problem we will consider three auxiliary problems

Problem 2 To find the real doubly-periodic function u^* harmonic in S^+ continuous in $S^- + L$, symmetric with respect to the axis $y = i\omega_2$ and satisfying the boundary condition

$$u^* = f_0(t), \quad t \in L, \quad (2)$$

where $f_0(t)$ is real symmetric doubly-periodic function given on L satisfies the H condition on L_{00}

Problem 3 To find the real function U^* sub-harmonic in S^+ continuous in $S^- + L$, of the class $P_e(1)$, symmetric with respect to the axis $y = i\omega_2$ satisfying the boundary condition (1) and the condition $U^*(x, y) > 0$.

Problem 4 To find a function $\Psi(z) = u_1 + iv_1$ of $P_e(1)$ class holomorphic in S^+ and continuous from the left everywhere on L , satisfying the following condition $|\Psi(t)| = f(t) \quad t \in L$.

Let us begin with the Problem 2.

We will find u^* as the real part of the sectionally holomorphic doubly periodic function $\Phi(z)$.

We assume that the contour L_{00} is smooth and the angle between its tangent and some constant direction satisfies the H condition.

Let $z = w(z^*)$ be conformal mapping of the upper z^* plane on the upper half of the rectangle S_{00} with the following correspondence of points: to the points $\omega_1 + 2i\omega_2$, $\omega_1 + i\omega_2$ corresponds the points $0, \infty$ consequently.

This mapping is given by the Shvartz-Kristoffel [1,4,6] formula.

$$w(z^*) = C \int_0^{z^*} \frac{dt}{\sqrt{(1-t^2)(1-k^2t^2)}} + \omega_1 + i\omega_2,$$

C and k are definite constants.

Continue this mapping analytically we obtain the mapping of the fundamental parallelogram on z plane. To the line L_{00} in z plane corresponds the line L^* in z^* plane.

The inverse transformation $z^* = w^{-1}(z)$ is an elliptic function with zeros at $\omega_1, 3\omega_1 + 2i\omega_2$ and with poles at $\omega_1 + i\omega_2, 2\omega_1 + 3i\omega_2$. The periods of this function are $4\omega_1$ and $2i\omega_2$ [1,4,6] and

$$z^* = C^* \frac{\sigma(z - \omega_1)\sigma(z - 3\omega_1 - 2i\omega_2)}{\sigma(z - \omega_1 - i\omega_2)\sigma(z - 2\omega_1 - 3i\omega_2)},$$

where σ is the Weierstrass "sigma-function" [1,4,9], C^* is the definite constant.

It is known from the theory of conformal mapping that, under the assumed conditions the functions, $w(z^*)$, $w^{-1}(z)$, are continuous from the left on L^* and L_{00} respectively.

It is obvious that if $f_0(t)$ satisfy the H condition on L_{00} , then the corresponding function obtained by the mapping $w^{-1}(z)$ will satisfy the H condition on L^* with the same Holder index.

So mapping the region S_{00} conformally on the plane and taking into the account that the function $\Phi(z)$ is doubly periodic and symmetric one arrives at exactly the same problem, but for the finite region S^* bounded by the line L^* , the boundary condition for the latter problem is given by the same formula if $f_0(t)$ is understood to the function $f_0(w(\tau))$ of points on the contour L^* .

Hence conclusions obtained for the case of region S^* may be directly transferred to the case of region S^- .

Using results of Muskhelishvili [8] the solution of the problem 2 in the plane z^* will be given by

$$u^*(z^*) = \operatorname{Re} \Phi^*(z^*),$$

where

$$\Phi^*(z^*) = \frac{1}{\pi i} \int_{L_{00}} \frac{\varphi(t)}{t - z^*} dt + iC_0, \quad (3)$$

$\varphi(t)$ is unknown real function satisfying the integral equation

$$\varphi(t_0) + \operatorname{Re} \left\{ \frac{1}{\pi i} \int_{L_{00}} \frac{\varphi(t)}{t - t_0} dt \right\} = f_0(t_0) + C_0, \quad t_0 \in L_{00}, \quad (4)$$

C_0 is the definite real constant.

This integral equation always has the unique solution [8]. So the initial problem is solvable.

Using the representation (3) and the transformation $t^* = w^{-1}(t)$ we obtain

$$\Phi(z) = \frac{1}{\pi i} \int_{L_{00}} \frac{\varphi(t)[w^{-1}(t)]' dt}{w^{-1}(t) - w^{-1}(z)} + iC_0, \quad (5)$$

Taking into the account symmetricity of the line L_{00} and of the function $\Phi^*(z^*)$ [8] and doubly-periodicity of the function

$$\frac{[w^{-1}(t)]' dt}{w^{-1}(t) - w^{-1}(z)}$$

we obtain the representation of the function $\Phi(z)$ in z plane in the following form

$$\Phi(z) = \frac{1}{\pi i} \int_{L_{00}} \varphi(t) [\zeta(t - z + \omega_1 + i\omega_2) + \zeta(t + z - \omega_1 - i\omega_2)] dt + iC_0, \quad (6)$$

where $\zeta(t - z)$ is the Weierstrass “zeta-function” and the integral equation (4) transforms to the following integral equation

$$\begin{aligned} \varphi(t_0) + \operatorname{Re} \left\{ \frac{1}{\pi i} \int_{L_{00}} \varphi(t) [\zeta(t - t_0 + \omega_1 + i\omega_2) + \zeta(t + t_0 - \omega_1 - i\omega_2)] dt \right\} \\ = f_0(t_0) + C_0, \quad t_0 \in L_{00}. \end{aligned} \quad (7)$$

using the results of Muskhelishvili [8] we can conclude that the integral equation (7) always has the solution. Hence, the solution of the Problem 2 is uniquely determined.

Now we consider the Problem 4. In S^+ let us consider the function

$$\Psi^*(z) = \ln \Psi(z),$$

where \ln is the definite branch of this function. The function $\Psi^*(z)$ is holomorphic in S^+ , of the class $\mathcal{P}(1)$ and satisfies the following boundary condition

$$\operatorname{Re} \Psi^*(t) = \ln f(t), \quad t \in L. \quad (8)$$

According to the previous results and results of the author [5] we obtain the representation of the function $\Psi^*(z)$ in the form

$$\begin{aligned} \Psi^*(z) = & \frac{1}{\pi i} \int_{L_{00}} \varphi(t) [\zeta(t - z + \omega_1 + i\omega_2) + \zeta(t + z - \omega_1 - i\omega_2)] dt \\ & + A \ln \sigma(z - \omega_1 - i\omega_2) + A \overline{\ln \sigma(z - \omega_1 - i\omega_2)} \\ & + D\zeta(z - \omega_1 - i\omega_2) + D \overline{\zeta(z - \omega_1 - i\omega_2)} \\ & + B(z - \omega_1 - i\omega_2)^2 + C_1(z - \omega_1 - i\omega_2) + iC_0, \end{aligned} \tag{9}$$

where $\zeta(t - z)$ is the Weierstrass “zeta-function”, C_0 is the definite real constant,

$$\begin{aligned} A &= \frac{\alpha_1 i\omega_2 - \alpha_2 \omega_1}{2\pi i}, \\ B &= \frac{\alpha_2 \delta_1 - \alpha_1 \delta_2}{4\pi i}, \\ D &= \frac{\beta_1 i\omega_2 - \beta_2 \omega_1 - (\alpha_1 - \alpha_2) i\omega_2 \omega_1}{2\pi i}, \\ C_1 &= \frac{\beta_2 \delta_1 - \beta_1 \delta_2 - i\omega_2 \delta_1 \alpha_2 + \omega_1 \delta_2 \alpha_1}{\pi i}, \end{aligned}$$

$\varphi(t)$ is unknown real function of H class, satisfying the integral equation

$$\begin{aligned} \varphi(t_0) + \operatorname{Re} \left\{ \frac{1}{\pi i} \int_{L_{00}} \varphi(t) [\zeta(t - t_0 + \omega_1 + i\omega_2) + \zeta(t + t_0 - \omega_1 - i\omega_2)] dt \right\} \\ = \ln f^*(t_0) + C_0, \quad t_0 \in L, \end{aligned} \tag{10}$$

where

$$\begin{aligned} f^*(t_0) = & f(t) - A \ln \sigma(z - \omega_1 - i\omega_2) - A \overline{\ln \sigma(z - \omega_1 - i\omega_2)} \\ & - D\zeta(z - \omega_1 - i\omega_2) - D \overline{\zeta(z - \omega_1 - i\omega_2)} \\ & - B(z - \omega_1 - i\omega_2)^2 - C_1(z - \omega_1 - i\omega_2). \end{aligned} \tag{11}$$

This integral equation has the unique solution. Finding $\varphi(t)$ the solution of the Problem 4 will be given by

$$\Psi(z) = \exp \Psi^*(z),$$

and the solution of the Problem 3 will be given by $U^*(z) = |\Psi(z)|$.

By means of the Problem 3 we now solve the initial problem .

For simplicity we assume that the function $u(x, y)$ has simple zeros at the points $a_1 + ia_2, 2i\omega_2 + a_1 - ia_2$.

Consider the function

$$\Phi^*(z) = \ln \Phi_1(z) + \ln \sigma(z - a_1 - ia_2) \sigma(z - 2i\omega_2 - a_1 + ia_2), \quad (12)$$

where $\Phi_1(z)$ is holomorphic function of $P_e(1)$ class with the same zeros as $u(z)$, subject the condition

$$u(z) = |\Phi_1(z)|. \quad (13)$$

The function $\Phi^*(z)$ is of the class $P_e(1)$ with the proper polynomials

$$\begin{aligned} P(z) &= \operatorname{Re}(\alpha_1 + 4\delta_1)z + \beta_1 - 2\delta_1(2a_1 - 2\omega_1 + i\omega_2), \\ Q(z) &= \operatorname{Re}(\alpha_2 + 4\delta_2)z + \beta_2 - 4\delta_1(a_1 - i\omega_2), \end{aligned}$$

According to the formulas (9), (12), (13) the solution of the Problem 1 is given by

$$u = |\Phi_1(z)| = \left| \frac{\exp \Phi^*(z)}{\sigma(z - a_1 - ia_2) \sigma(z - 2i\omega_2 - a_1 + ia_2)} \right|,$$

where $\Phi^*(z)$ is the solution of the Problem 4.

Let us check the uniqueness of the Problem 1.

Let u_1 and u_2 be two possible solutions of the Problem 1.

The function

$$\Phi_0(z) = \ln \Phi_1(z) - \ln \Phi_2(z),$$

where $u_1 = |\Phi_1(z)|$, $u_2 = |\Phi_2(z)|$, will be doubly-periodic with no poles, satisfying the condition $\operatorname{Re} \Phi_0(z) = 0$. Hence $\Phi_0(z) = C$, C is the constant [1,4,6,8].

In some particular cases the solution of the Problem 3 and hence of the Problem 1 can be written effectively:

1) In the case of L_{00} is a rectangle $ABCD$ (Fig.1), we immediately obtain

$$\begin{aligned} U^*(z) &= \left| \exp \left\{ \frac{1}{\pi i} \int_{BC \cup AD} f^*(t) [\zeta(t - t_0) + \zeta(t + t_0)] dt \right. \right. \\ &\quad + \frac{1}{\pi} \int_{AB \cup CD} f^*(t) [\zeta(it - it_0) + \zeta(it + it_0)] dt \\ &\quad + A \ln \sigma(z - \omega_1 - i\omega_2) + A \overline{\ln \sigma(z - \omega_1 - i\omega_2)} \\ &\quad + D \zeta(z - \omega_1 - i\omega_2) + D \overline{\zeta(z - \omega_1 - i\omega_2)} \\ &\quad \left. + B(z - \omega_1 - i\omega_2)^2 + C_1(z - \omega_1 - i\omega_2) + iC_0 \right|, \end{aligned}$$

where BC and AD are horizontal sides of the rectangle.

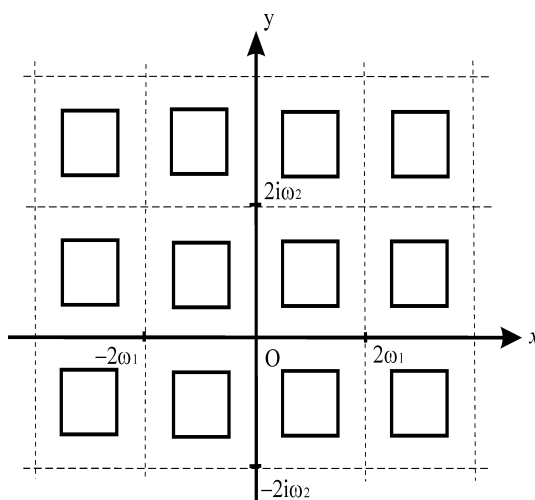


Figure 1: The case of rectangle.

2) In the case of L_{00} is the circle with the sufficiently small radius $r = \varepsilon$ and with the center (ω_1, ω_2) (Fig. 2). Using the behavior of conformal mapping and the Shvartz formula [5,8] we obtain

$$\begin{aligned}
 U^*(z) = & \left| \exp \left\{ \frac{1}{\pi i} \int_{L_{00}} f^*(t) [\zeta(t - t_0) + \zeta(t + t_0)] dt \right. \right. \\
 & + A \ln \sigma(z - \omega_1 - i\omega_2) + \overline{A \ln \sigma(z - \omega_1 - i\omega_2)} \\
 & + D\zeta(z - \omega_1 - i\omega_2) + \overline{D\zeta(z - \omega_1 - i\omega_2)} \\
 & \left. \left. + B(z - \omega_1 - i\omega_2)^2 + C_1(z - \omega_1 - i\omega_2) + iC_0 \right\} \right|.
 \end{aligned}$$

Note. The doubly-periodic problems for the doubly-periodic areas was considered in [3] and [10].

Note. In the case of $\alpha_1 = \alpha_2 = 0, \beta_1 = \beta_2 = \beta$, where β is a real constant, the function $\text{Re } \Psi(z)$ will be harmonic of the class $P_e(0)$

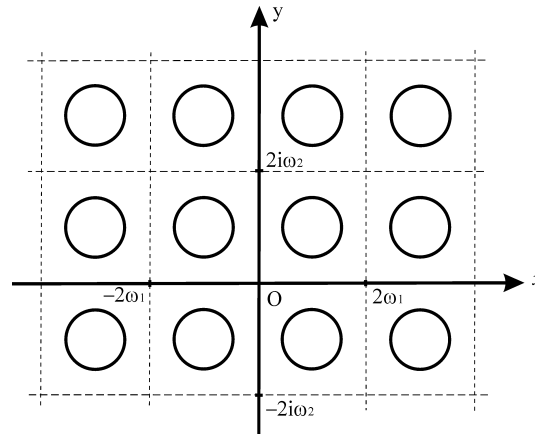


Fig.1. The case of $\text{Im } 2\omega_1=0$; $\text{Re } 2i\omega_2=0$.
 L_{00} is the circle.

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